

A MASS TRANSPORT APPROACH FOR THE RELATIVISTIC HEAT EQUATION

Marjolaine Puel

Institut de Mathématiques de Toulouse

Joint work with R. McCann and also X. Cabre and J. A. Carrillo

email : puel@mip.ups-tlse.fr

The relativistic heat equation

We are interested in the Cauchy problem for the relativistic heat equation

$$\partial_t \rho = \operatorname{div} \left(\rho \frac{\nabla \rho}{\sqrt{\rho^2 + |\nabla \rho|^2}} \right) = \operatorname{div} \left(\rho \frac{\nabla \log \rho}{\sqrt{1 + |\nabla \log \rho|^2}} \right). \quad (1)$$

introduced to impose an upperbound for the propagation velocity.

(Ref Brenier (01), Rosenau (92), Mihalas-Mihalas (84))

Assumptions $\rho(t, x)$ with $t \in [0, T]$ and $x \in \Omega$, bounded domain of \mathbf{R}^d and

$$0 < m \leq \rho_0 \leq M \text{ and } \int_{\Omega} \rho dx = 1.$$

This problem has been independently adressed in

Andreu, Caselles, Mazon (ref Non Linear Analysis 2005 and JEMS 2005)

Mass Transport Strategy

The aim of our work is to implement a different point of view following the ideas of [Jordan, Kinderlehrer, Otto SIAM J.Math. Anal.\(98\)](#)

→ construction of solutions of general equation of the type

$$\partial_t \rho = \operatorname{div} (\rho \nabla c^* (\nabla (F'(\rho))))$$

where F is a [convex Entropy](#) and c is a [convex cost function](#)

c^* it's Legendre transform i.e. $c^*(x) = \sup_{y \in \mathbf{R}^d} x \cdot y - c(y)$.

→ time discrete scheme involving a [double minimization process](#).

Time discrete scheme

Let $P(\Omega)$ be the set of probability measures on Ω ,

$\rho_0 \in P(\Omega)$ given, find $\rho^h(t, x) \in P(\Omega)$ defined by

$$\begin{cases} \rho^h(0, x) = \rho_0(x) \\ \rho^h(t, x) = \rho_i^h(x) \quad \text{for } t \in]ih; (i+1)h] \quad h \text{ being the time step} \end{cases} \quad (2)$$

where

$$\rho_i^h(x) = \operatorname{argmin}_{\rho \in P(\Omega)} \left(\int_{\Omega} F(\rho(y)) dy + h \inf_{\gamma \in \Gamma_i^h(\rho_{i-1}^h, \rho)} \int_{\Omega \times \Omega} c\left(\frac{x-y}{h}\right) d\gamma(x, y) \right),$$

$\Gamma_i^h(\rho_{i-1}^h, \rho)$ is the set of transport plans between ρ_{i-1}^h and ρ .

→ Generalisation of **discrete gradient flow** to general convex cost function

Ref **Ambrosio Gigli Savaré Lectures in math ETH (2005)**,

Villani Graduate studies in Math AMS (2003)

Previous results

- **Jordan-Kinderlehrer-Otto SIAM Journ Math Anal (98)**

$$c(z) = \frac{|z|^2}{2} \text{ and } F(\rho) = \beta^{-1} \rho \log \rho - V \rho, V \text{ given.}$$

→ Linear Fokker Plank equation

$$\partial_t \rho = \operatorname{div} (\nabla V \rho) + \beta^{-1} \Delta \rho.$$

- **Otto Preprint (96)** $c(z) = \frac{|z|^q}{q}$ and $F(\rho) = \frac{n \rho^m}{m(m-1)}$ where
 $m = n + \frac{p-2}{p-1}, \frac{1}{p} + \frac{1}{q} = 1, p \geq 2$

→ Doubly degenerate equation

$$\partial_t \rho = \operatorname{div} (|\nabla \rho^n|^{p-2} \nabla \rho^n)$$

- **Agueh Adv Diff Equ (05)** $\beta |z|^q \leq c(z) \leq \alpha(|z|^q + 1)$ and F convex + displacement convexity

→ A large set of equation

$$\partial_t \rho = \operatorname{div} (\rho \nabla c^* (\nabla (F'(\rho))))$$

Cost Function and Entropy for the relativistic heat equation

The relativistic heat equation does not belong to those set of equations since

$$c(z) = \begin{cases} 1 - \sqrt{1 - |z|^2} & \text{if } |z| \leq 1 \\ \infty & \text{elsewhere .} \end{cases}$$

and $F(\rho) = \rho \log \rho - \rho$.

This cost function is strictly convex and discontinuous ($c(z) = \infty$ if $|z| > 1$)

\iff the velocity ∇c^* is bounded

\implies characteristic property of a relativistic phenomenon.

General assumptions on the cost and on the Entropy

This present work will in fact apply for any

- Cost function

$$c(z) = \begin{cases} \tilde{c}(|z|) \geq 0 \text{ if } |z| \leq 1 \\ \infty \text{ elsewhere} \end{cases}$$

where \tilde{c} is strictly convex, $C^0[0, 1] \cup C^2([0, 1[)$,

with $|\nabla c(z)| \rightarrow \infty$ when $|z| \rightarrow 1$.

- Entropy function $F \in C^2(\mathbf{R})$ satisfying $\frac{F(\lambda)}{\lambda} \rightarrow \infty$ when $\lambda \rightarrow \infty$

and $\lambda^d F(\lambda^{-d})$ is convex \rightarrow displacement

convexity assumption,

see McCann Adv. Math. (97)

Formal argument I

Step 1 Find γ_i^h for each time intervall $[ih; (i+1)h]$, its second marginal ρ_i^h and S_i^h , the associated optimal map i.e. $\gamma_i^h = \delta(x - S_i^h(y))$.

The existence of ρ_i^h is ensured by the double minimization process.

Problem 1 Definition of the map

Step 2 Derive the Euler-Lagrange equations (dervation / ρ and then ∇)

$$\nabla(F'(\rho(y))) = \nabla c\left(\frac{S_i^h(y) - y}{h}\right) \implies \frac{S_i^h(y) - y}{h} = \nabla c^*(\nabla(F'(\rho(y))))$$

since $\nabla c^*(\nabla c(z)) = z$.

Problem 2 We prove $\{|S_i^h(y) - y| \leq h\}$ but if $|S_i^h(y) - y| = h$,

then $\nabla c\left(\frac{S_i^h(y) - y}{h}\right)$ is not defined.

Formal argument II

Step 3 Obtain an approximate time discrete equation for ρ^h by $\times \rho \nabla \phi$

$$\frac{\rho_i^h - \rho_{i-1}^h}{h} = \operatorname{div}(\rho_i^h \nabla c^*(\nabla(F'(\rho_i^h)))) + \text{Correction terms in } O(h).$$

Step 4 It remains to pass to the limit when the time step h goes to zero

$$\implies \partial_t \rho = \operatorname{div}(\rho \nabla c^*(\nabla(F'(\rho)))).$$

Problem 3 We want to use a monotonicity argument (Minty-Browder) to identify the limit but **lack of regularity**.

Construction of the optimal map (drop the h)

The classical result of Gangbo McCann CRAS (95) can not be applied here. Indeed it is based on the Kantorovich duality:

$$\int_{\Omega \times \Omega} c(x - y) d\gamma_{opt}(x, y) = \int_{\Omega} \phi(x) d\rho_0(x) + \int_{\Omega} \psi(y) d\rho_1(y)$$

where ϕ is the c -transform of ψ , i.e.

$$\phi(x) = \inf_y (c(x - y) - \psi(y))$$

When $c \in C^{1,\alpha}$, ψ is Lipschitz and then differentiable which means that when $(x, y) \in \text{supp}\gamma_{opt}$, we have

$$\nabla c(x - y) = -\nabla \psi(y) \text{ and then } x = y + \nabla c^*(-\nabla \psi(y))$$

In our case, ψ is no more Lipschitz, we have a problem when $|x - y| = 1$.

The strategy consists in introducing a mollified problem.

The mollified problem

We introduce the [Yoshida mollification](#) of the convex function c

[Ref Brezis Operateur maximaux monotones et semi-groupes de contraction dans les espaces de Hilbert \(73\)](#) by a

Convexification of c^*

$$c^{\epsilon*}(x) = c^*(x) + \frac{\epsilon}{2}|x|^2$$

\implies Mollification of c

$$c^\epsilon(x) = \inf_{z \in \mathbf{R}^d} \left(c(x - z) + \frac{|z|^2}{2\epsilon} \right)$$

and we define $\gamma_i^{\epsilon h}$, the optimal transport plan between ρ_{i-1}^h and $\rho_i^{\epsilon h}$ obtained when the minimization involves the mollified cost function c^ϵ .

Kantorovich duality and optimal map for the mollified problem

$$\int_{\Omega \times \Omega} c^\epsilon(x - y) d\gamma^\epsilon(x, y) = \int_{\Omega} \phi^\epsilon(x) d\rho_0(x) + \int_{\Omega} \psi^\epsilon(y) d\rho^\epsilon(y)$$

where the potential function satisfies

$$\phi^\epsilon(x) = \inf_{y \in \mathbf{R}^d} (c^\epsilon(x - y) - \psi^\epsilon(y)).$$

Gangbo-McCann \Rightarrow existence of a map $S^\epsilon(y) = y + \nabla c^{\epsilon*}(-\nabla \psi^\epsilon(y))$

Agueh \Rightarrow Euler-Lagrange eq. $S^\epsilon(y) = y + \nabla c^{\epsilon*}(\nabla(F'(\rho^\epsilon(y))))$

and then $\psi^\epsilon(y) = -F'(\rho^\epsilon(y))$ and $\rho^\epsilon(y)$ is bounded in $W^{1,1}(\Omega)$

since we have the entropy-information inequality

$$\int \rho^\epsilon \nabla c^{\epsilon*}(\nabla(F'(\rho^\epsilon))) \cdot \nabla(F'(\rho^\epsilon)) dy \leq \int [F(\rho_0(y)) - F(\rho^\epsilon(y))] dy.$$

Limiting process to obtain the optimal map

1- the limit $\lim_{\epsilon \rightarrow 0} \gamma_i^{\epsilon h} = \gamma_i^h$ and

$\text{supp } \gamma_i^h \subset \{(x, y) \text{ such that } |x - y| \leq h\} \cap Z_i^h$ where $\gamma_i^h(Z_i^h) = 0$.

2- $\lim_{\epsilon \rightarrow 0} \rho_i^{\epsilon h} = \rho_i^h \in BV(\Omega) \implies$ approximatively differentiable

\implies avoid a.e. the undetermination of the map.

Ref Ambrosio Gigli Savaré Lectures in math ETH (2005),

3- Kantorovich duality

$$\int_{\Omega \times \Omega} hc\left(\frac{x-y}{h}\right) d\gamma_i^h(x, y) = \int_{\Omega} \phi_i^h(x) d\rho_{i-1}^h(x) + \int_{\Omega} \psi_i^h(y) d\rho_i^h(y).$$

with $\phi_i^h(x) = \inf(y) (hc(\frac{x-y}{h}) - \psi_i^h(y))$ and $\psi_i^h(y) = -F'(\rho_i^h)$

\implies **Optimal map** + **Euler-Lagrange equation**

$$S_i^h(y) = y + h\nabla c^*(-\nabla\psi_i^h(y)) = y + h\nabla c^*(\nabla(F'(\rho_i^h(y)))).$$

From the time discrete equation to the continuous equation

By multiplying the Euler-Lagrange equation by $\rho_i^h \nabla \phi$, we obtain the approximate time discrete equation

$$\frac{\rho_i^h - \rho_{i-1}^h}{h} = \operatorname{div}(\rho_i^h \nabla c^*(\nabla(F'(\rho_i^h)))) + \text{Correction terms in } O(h).$$

and when h goes to zero, we obtain the continuous equation

$$\partial_t \rho = \operatorname{div}(\rho A)$$

where ρ is the $L^1([0, T] \times \Omega)$ limit of ρ^h

and A is the $w^*L^\infty(\Omega)$ limit of $\nabla c^*(\nabla(F'(\rho^h)))$.

It remains to identify the limit A

Main Result

Theorem

(i) **Support of the optimal measure: Finite speed of propagation**

$$\text{supp } \gamma_i^h \subset \{(x, y) \mid \frac{|x - y|}{h} < 1\} \cup Z_i^h \text{ with } \gamma_i^h(Z_i^h) = 0.$$

(ii) **Euler-Lagrange equation: a discrete scheme**

$$\gamma_i^h(x, y) = \delta(x - S_i^h(y)) \text{ with } S_i^h(y) = y + h \nabla c^*(\nabla(F'(\rho_i^h(y)))).$$

(iii) **Convergence of the measure ρ^h .** Up to a subsequence,

$$\rho^h \longrightarrow \bar{\rho} \text{ in } L^1([0, T] \times \Omega) \text{ and w-}L_w^1([0, T], BV(\Omega))$$

$$\text{and } \bar{\rho} \in L^\infty([0, T] \times \Omega) \cap L_w^1([0, T], BV(\Omega)).$$

(iv) **Limiting equation $\bar{\rho}$** is a solution in the sense of distribution to

$$\partial_t \bar{\rho} = \text{div}(\bar{\rho} \nabla c^*(\nabla(F'(\bar{\rho}))).$$