

Hamiltonian Reductions Relating To Optimal Mass Transport

Paul Lee

Department of Mathematics

University of Toronto

Email: plee@math.toronto.edu

(Joint Work with Boris Khesin)

The Diffeomorphism Group

Let (M, \langle, \rangle) be a Riemannian compact manifold without boundary.

The Diffeomorphism Group

Let (M, \langle, \rangle) be a Riemannian compact manifold without boundary.

- Diffeomorphism Group $Diff$ is the space of all diffeomorphisms on M .

The Diffeomorphism Group

Let (M, \langle, \rangle) be a Riemannian compact manifold without boundary.

- Diffeomorphism Group $Diff$ is the space of all diffeomorphisms on M .
- Infinite-dimensional Lie group with Lie algebra given by the space of all vector fields, denoted by $Vect$.

The Diffeomorphism Group

Let (M, \langle, \rangle) be a Riemannian compact manifold without boundary.

- Diffeomorphism Group $Diff$ is the space of all diffeomorphisms on M .
- Infinite-dimensional Lie group with Lie algebra given by the space of all vector fields, denoted by $Vect$.
- Tangent space $T_\phi Diff$ at $\phi \in Diff$ is given by

$$\{X \circ \phi : M \rightarrow TM \mid X \text{ is a vector field}\}.$$

Symmetry Group

- Fix a volume form μ , consider the subgroup of all diffeomorphisms which preserve μ . We denote this subgroup by $SDiff$.

Symmetry Group

- Fix a volume form μ , consider the subgroup of all diffeomorphisms which preserve μ . We denote this subgroup by $SDiff$.
- The Lie algebra of $SDiff$ is the space of all vector fields divergence free with respect to the volume form μ , denoted by $SVect$.

Symmetry Group

- Fix a volume form μ , consider the subgroup of all diffeomorphisms which preserve μ . We denote this subgroup by $SDiff$.
- The Lie algebra of $SDiff$ is the space of all vector fields divergence free with respect to the volume form μ , denoted by $SVect$.
- $SDiff$ acts on the diffeomorphism group

$$\Phi_g(\phi) = \phi \circ g, \quad g \in SDiff, \quad \phi \in Diff.$$

Moment Map

- The pushforward of the action Φ_g defined an action on $TDiff$ which is given by

$$T\Phi_g(X \circ \phi) = X \circ \phi \circ g, \quad X \circ \phi \in TDiff, \quad g \in SDiff.$$

Moment Map

- The pushforward of the action Φ_g defined an action on $TDiff$ which is given by

$$T\Phi_g(X \circ \phi) = X \circ \phi \circ g, \quad X \circ \phi \in TDiff, \quad g \in SDiff.$$

- The corresponding moment map $\mu : TDiff \rightarrow SVect^*$ is given by

$$\mu(X \circ \phi)(\xi) = \int_M \langle X, \phi_* \xi \rangle \phi_* \mu, \quad \xi \in SVect.$$

Burgers Equation as Hamiltonian System

- Define a Riemannian metric on the diffeomorphism group $Diff$:

$$\langle X_1 \circ \phi, X_2 \circ \phi \rangle^{Diff} = \int_M \langle X_1 \circ \phi, X_2 \circ \phi \rangle \mu.$$

Burgers Equation as Hamiltonian System

- Define a Riemannian metric on the diffeomorphism group $Diff$:

$$\langle X_1 \circ \phi, X_2 \circ \phi \rangle^{Diff} = \int_M \langle X_1 \circ \phi, X_2 \circ \phi \rangle \mu.$$

- If $H^{Diff}(X \circ \phi) = \frac{1}{2} \langle X \circ \phi, X \circ \phi \rangle^{Diff}$ is the corresponding kinetic energy Hamiltonian and ϕ_t is the corresponding geodesic, then it satisfies the Burgers equation:

$$\partial_t \phi = X_t \circ \phi_t, \quad \partial_t X + \nabla_{X_t} X_t = 0.$$

Solution of Burgers Equation

- Geodesics on $Diff$ starting at id are of the form $\phi_t : x \mapsto \exp(tX(x))$, where \exp is the Riemannian exponential and X is the initial velocity field.

Solution of Burgers Equation

- Geodesics on $Diff$ starting at id are of the form $\phi_t : x \mapsto \exp(tX(x))$, where \exp is the Riemannian exponential and X is the initial velocity field.
- If X_t is the corresponding velocity field of the geodesic ϕ_t , then $X_t \circ \phi_t$ is a trajectory of the Hamiltonian flow of the kinetic energy.

Conservation of Momentum

Moment Map: $\mu(X \circ \phi)(\xi) = \int_M \langle X, \phi_* \xi \rangle \phi_* \mu.$
 $\mu : TDiff \rightarrow SVect^*$

Geodesics starting at id : $x \mapsto \exp(tX(x)).$

Conservation of Momentum

Moment Map: $\mu(X \circ \phi)(\xi) = \int_M \langle X, \phi_* \xi \rangle \phi_* \mu.$
 $\mu : TDiff \rightarrow SVect^*$

Geodesics starting at id : $x \mapsto \exp(tX(x)).$

- **By Hodge decomposition,** $\mu^{-1}(0) = \{\nabla f \circ \phi \mid f \in C^\infty(M)\}.$

Conservation of Momentum

Moment Map: $\mu(X \circ \phi)(\xi) = \int_M \langle X, \phi_* \xi \rangle \phi_* \mu.$,
 $\mu : TDiff \rightarrow SVect^*$

Geodesics starting at id : $x \mapsto \exp(tX(x)).$

- **By Hodge decomposition,** $\mu^{-1}(0) = \{\nabla f \circ \phi \mid f \in C^\infty(M)\}.$
- **Conservation of Momentum:** For each fixed $z \in SVect$, the Hamiltonian flow of H^{Diff} preserve the level set $\mu^{-1}(z).$

Conservation of Momentum

Moment Map: $\mu(X \circ \phi)(\xi) = \int_M \langle X, \phi_* \xi \rangle \phi_* \mu.$
 $\mu : TDiff \rightarrow SVect^*$

Geodesics starting at id : $x \mapsto \exp(tX(x)).$

- **By Hodge decomposition,** $\mu^{-1}(0) = \{\nabla f \circ \phi \mid f \in C^\infty(M)\}.$
- **Conservation of Momentum:** For each fixed $z \in SVect$, the Hamiltonian flow of H^{Diff} preserve the level set $\mu^{-1}(z).$
- **Consider $z = 0$,** the solutions of the Burgers equation with gradient initial velocity field $\phi_t = \exp(t\nabla f)$ satisfies $\partial_t \phi_t = \nabla f_t \circ \phi_t$, where $f_t \in C^\infty(M)$ for each $t.$

Other Hamiltonian Systems

- Let $H : TM \rightarrow \mathbb{R}$ and let H^{Diff} be the Hamiltonian on $TDiff$

$$H^{Diff}(X \circ \phi) = \int_M H(X \circ \phi) \mu.$$

Other Hamiltonian Systems

- Let $H : TM \rightarrow \mathbb{R}$ and let H^{Diff} be the Hamiltonian on $TDiff$

$$H^{Diff}(X \circ \phi) = \int_M H(X \circ \phi) \mu.$$

- **Conservation of Momentum:** Let Φ_t^H be the Hamiltonian flow of H on TM and let $\pi : TM \rightarrow M$ be the tangent bundle projection. Set $\phi_t = \pi \circ \Phi_t^H \circ \nabla f$, then $\Phi_t^H = \nabla f_t \circ \phi_t$, where f_t are functions on the manifold M .

Other Hamiltonian Systems

- Let $H : TM \rightarrow \mathbb{R}$ and let H^{Diff} be the Hamiltonian on $TDiff$

$$H^{Diff}(X \circ \phi) = \int_M H(X \circ \phi) \mu.$$

- **Conservation of Momentum:** Let Φ_t^H be the Hamiltonian flow of H on TM and let $\pi : TM \rightarrow M$ be the tangent bundle projection. Set $\phi_t = \pi \circ \Phi_t^H \circ \nabla f$, then $\Phi_t^H = \nabla f_t \circ \phi_t$, where f_t are functions on the manifold M .
- **Method of Characteristics:** f_t above satisfies the Hamilton-Jacobi equation $\partial_t f_t + H(\nabla f_t) = 0$.

Reduction of Dynamics

- The quotient $Diff/SDiff$ can be identified with the space \mathcal{W} of all volume forms with total integral equal to that of μ ,
 $\Pi : \phi \mapsto \phi_*\mu$.

Reduction of Dynamics

- The quotient $Diff/SDiff$ can be identified with the space \mathcal{W} of all volume forms with total integral equal to that of μ ,
 $\Pi : \phi \mapsto \phi_*\mu$.
- Marsden Weinstein quotient:

$$\mu^{-1}(0)/SDiff = T(Diff/SDiff) \cong T\mathcal{W}.$$

Reduction of Dynamics(cont')

- The kinetic energy $H^{Diff}(X \circ \phi) = \int_M \langle X_1 \circ \phi, X_2 \circ \phi \rangle \mu$ is invariant under the $SDiff$ -action on $TDiff$. It induces a kinetic energy Hamiltonian $H^{\mathcal{W}}$ on $T\mathcal{W}$.

Reduction of Dynamics(cont')

- The kinetic energy $H^{Diff}(X \circ \phi) = \int_M \langle X_1 \circ \phi, X_2 \circ \phi \rangle \mu$ is invariant under the $SDiff$ -action on $TDiff$. It induces a kinetic energy Hamiltonian $H^{\mathcal{W}}$ on $T\mathcal{W}$.
- By conservation of momentum, the Hamiltonian flow of H^{Diff} restricts to a flow on the the level $\mu^{-1}(0)$. This flow can also be reduced to a flow on the Marsden-Weinstein quotient $T\mathcal{W}$.

Reduction of Dynamics(cont')

- The kinetic energy $H^{Diff}(X \circ \phi) = \int_M \langle X_1 \circ \phi, X_2 \circ \phi \rangle \mu$ is invariant under the $SDiff$ -action on $TDiff$. It induces a kinetic energy Hamiltonian $H^{\mathcal{W}}$ on $T\mathcal{W}$.
- By conservation of momentum, the Hamiltonian flow of H^{Diff} restricts to a flow on the the level $\mu^{-1}(0)$. This flow can also be reduced to a flow on the Marsden-Weinstein quotient $T\mathcal{W}$.
- Reduction of Dynamics: The Hamiltonian flow of $H^{\mathcal{W}}$ is the reduced flow described above.

Reduction of Dynamics(cont')

- The kinetic energy $H^{Diff}(X \circ \phi) = \int_M \langle X_1 \circ \phi, X_2 \circ \phi \rangle \mu$ is invariant under the $SDiff$ -action on $TDiff$. It induces a kinetic energy Hamiltonian $H^{\mathcal{W}}$ on $T\mathcal{W}$.
- By conservation of momentum, the Hamiltonian flow of H^{Diff} restricts to a flow on the the level $\mu^{-1}(0)$. This flow can also be reduced to a flow on the Marsden-Weinstein quotient $T\mathcal{W}$.
- Reduction of Dynamics: The Hamiltonian flow of $H^{\mathcal{W}}$ is the reduced flow described above.
- The geodesics on \mathcal{W} corresponding to the Hamiltonian $H^{\mathcal{W}}$ are $\exp(t\nabla f)_*\mu$.

Relations with Optimal Transport

- **Optimal Transport:** Let d be the Riemannian distance function. Find $\phi : M \rightarrow M$ which minimizes

$$\int_M d^2(x, \phi(x)) \mu$$

among all maps satisfying $\phi_*\mu = \nu$.

Relations with Optimal Transport

- **Optimal Transport:** Let d be the Riemannian distance function. Find $\phi : M \rightarrow M$ which minimizes

$$\int_M d^2(x, \phi(x)) \mu$$

among all maps satisfying $\phi_*\mu = \nu$.

- (Brenier 91', McCann 94') Optimal maps are of the form $\exp(t\nabla f)$.

Relations with Optimal Transport

- **Optimal Transport:** Let d be the Riemannian distance function. Find $\phi : M \rightarrow M$ which minimizes

$$\int_M d^2(x, \phi(x)) \mu$$

among all maps satisfying $\phi_*\mu = \nu$.

- (Brenier 91', McCann 94') Optimal maps are of the form $\exp(t\nabla f)$.
- (Otto 00') The curve $\exp(t\nabla f)_*\mu$ are geodesics on the Wasserstein space.

Relations with Optimal Transport(cont')

- Let $L : TM \rightarrow \mathbb{R}$ be a Lagrangian and let $c : M \times M \rightarrow \mathbb{R}$ be the cost function defined by

$$c(x, y) = \inf_{\{\gamma: [0,1] \rightarrow M \mid \gamma(0)=x, \gamma(1)=y\}} \int_0^1 L(\gamma(t), \dot{\gamma}(t)) dt.$$

Relations with Optimal Transport(cont')

- Let $L : TM \rightarrow \mathbb{R}$ be a Lagrangian and let $c : M \times M \rightarrow \mathbb{R}$ be the cost function defined by

$$c(x, y) = \inf_{\{\gamma: [0,1] \rightarrow M \mid \gamma(0)=x, \gamma(1)=y\}} \int_0^1 L(\gamma(t), \dot{\gamma}(t)) dt.$$

- **Optimal Transport:** Find $\phi : M \rightarrow M$ such that $\phi_*\mu = \nu$ and achieves the infimum

$$\inf_{\{\phi \mid \phi_*\mu = \nu\}} \int_M c(x, \phi(x)) \mu.$$

Relations with Optimal Transport(cont')

- Let $L : TM \rightarrow \mathbb{R}$ be a Lagrangian and let $c : M \times M \rightarrow \mathbb{R}$ be the cost function defined by

$$c(x, y) = \inf_{\{\gamma: [0,1] \rightarrow M \mid \gamma(0)=x, \gamma(1)=y\}} \int_0^1 L(\gamma(t), \dot{\gamma}(t)) dt.$$

- **Optimal Transport:** Find $\phi : M \rightarrow M$ such that $\phi_*\mu = \nu$ and achieves the infimum

$$\inf_{\{\phi \mid \phi_*\mu = \nu\}} \int_M c(x, \phi(x)) \mu.$$

Relations with Optimal Transport(cont')

- Let $H(x, v) = \inf_{u \in T_x M} (\langle u, v \rangle - L(x, u))$ be the Legendre transform of the function L .

Relations with Optimal Transport(cont')

- Let $H(x, v) = \inf_{u \in T_x M} (\langle u, v \rangle - L(x, u))$ be the Legendre transform of the function L .
- (Bernard, Buffoni 04') Optimal maps are of the form

$$\pi \circ \Phi_t \circ \nabla f$$

where Φ_t is the Hamiltonian flow of the Hamiltonian H and $\pi : TM \rightarrow M$ is the tangent bundle projection.

Relations with Optimal Transport(cont')

- Let $H(x, v) = \inf_{u \in T_x M} (\langle u, v \rangle - L(x, u))$ be the Legendre transform of the function L .
- (Bernard, Buffoni 04') Optimal maps are of the form

$$\pi \circ \Phi_t \circ \nabla f$$

where Φ_t is the Hamiltonian flow of the Hamiltonian H and $\pi : TM \rightarrow M$ is the tangent bundle projection.

- We can see this as a consequence of Hamiltonian reduction.